

Hybrid Thermally Activated Nanocluster Fluorophores for X-ray Scintillators

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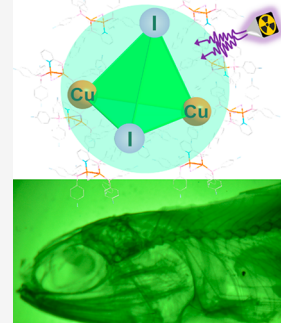
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ABSTRACT: Thermally activated delayed fluorescence (TADF) X-ray scintillators have attracted increasing attention because of their theoretical ability to utilize 100% radiation-induced excitons. However, their scintillation performance is severely hampered by their typically low X-ray attenuation efficiencies or mechanochromic properties. Here, we demonstrate hybrid organic–inorganic TADF scintillators with remarkably high X-ray absorption cross sections based on non-mechanochromic Cu_2X_2 ($\text{X} = \text{Cl}, \text{Br}, \text{or I}$) nanoclusters. The Cu_2X_2 nanoclusters display radioluminescence light yields as high as $175\,000$ photons MeV^{-1} , attributable to their low self-absorption and spatially separated HOMO and LUMO orbitals. Furthermore, because of their non-mechanochromic properties, these nanoclusters can be formed into scintillating screens with excellent radiation and humidity stability via soft-pressing. The fabricated TADF scintillating screens display a high X-ray imaging efficiency, benefiting from their high light yields and high spatial resolutions (~ 30 lp mm^{-1}). This work highlights the promise TADF Cu_2X_2 nanocluster-based scintillators have in meeting the demanding requirements of cutting-edge X-ray imaging.

High-Performance X-Ray Imaging with TADF Copper Halide Nanoclusters



X-ray scintillation plays an indispensable role in medical diagnosis, industrial inspection, and security, given its ability to convert high-energy ionizing radiation into readily detectable low-energy visible photons.^{1–5} New scintillators with enhanced abilities for converting X-rays to visible light are actively being explored to improve radioluminescence (RL) intensity and imaging quality.^{6–14} One promising avenue of investigation is the use of thermally activated delayed fluorescence (TADF) materials,^{9–14} particularly the hybrid organic–inorganic TADF-like Cu(I) complexes. Theoretically, these TADF-like Cu(I) complexes can achieve up to 100% internal quantum efficiency through the reverse intersystem crossing (RISC) process. This process facilitates the transition of their excitons from the lowest triplet excited state (T_1) to the lowest singlet excited state (S_1),^{15–17} thereby significantly enhancing the efficient utilization of radiation-induced triplet excitons.

Compared to organic TADF scintillators, TADF-like Cu(I) complexes are ideal systems for high-performance scintillators, as they contain heavy atoms and display appreciable RL properties under ambient conditions.^{9,11–13} The structural tunability of Cu(I) complexes allows for one to obtain a small energy gap (ΔE_{ST}) between S_1 and T_1 , as well as strong spin–orbit coupling (SOC).^{18–21} The combination of these properties opens up new possibilities for radiation detection and imaging,^{12,13} as it can overcome the trade-off between high

quantum yield (QY) and small ΔE_{ST} values that organic TADF systems have.^{22–24}

These advantages have spurred the development of scintillators in the family of hybrid organic–inorganic TADF Cu(I) halide complexes, such as Cu_4I_4 ,^{11,25} Cu_3Br_3 ,²⁶ and Cu_1I_1 ,⁹ which have shown promising light yields ($\sim 12\,842$, $\sim 20\,000$, and $\sim 30\,000$ photons MeV^{-1} , respectively) and low detection limits. The radioluminescence properties of these materials are attributed to the combination of their heavy [CuI/Br]_n core and high photoluminescence efficiency.^{9,11,25,26} While previous studies have confirmed the potential of Cu–I nanoclusters for scintillator applications,^{2,7,27,28} the light yields of such nanoclusters still require significant improvement in order to be on par with state-of-the-art imaging scintillators.^{7,8,10,29} Moreover, their stability and mechanochromism^{30,31} are likely to hinder their processing into high-quality scintillator screens via low-cost mechanical methods.^{32–34}

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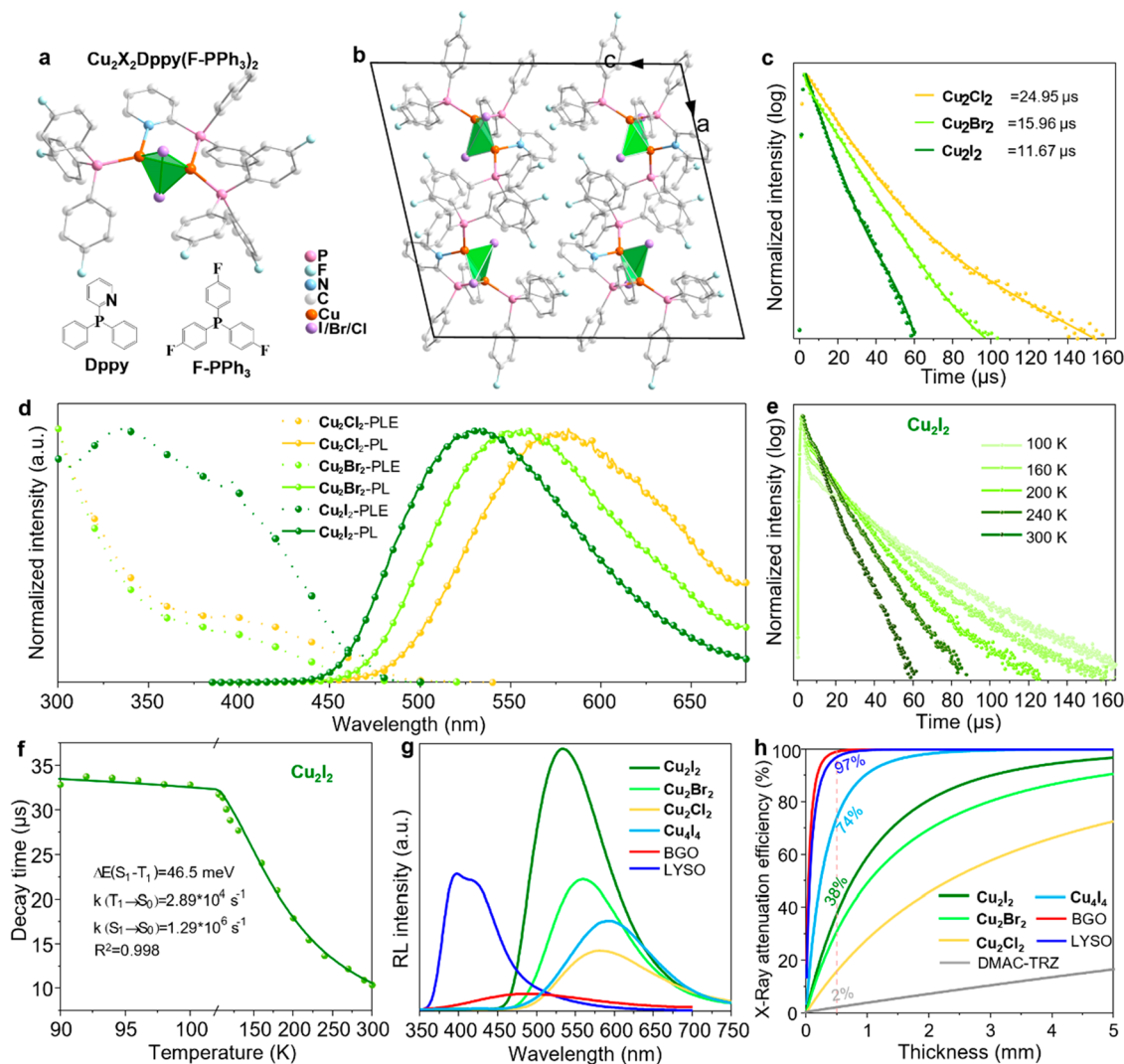


Figure 1. Characterization of the Cu_2X_2 crystals. (a) Molecular structures and (b) packing mode of the Cu_2X_2 crystals. (c) Excited-state lifetimes upon excitation at 365 nm. (d) Normalized emission and excitation spectra. (e) Temperature dependence of the excited-state lifetimes in the range of 100–300 K for Cu_2I_2 . (f) Plot of emission decay lifetime against temperature (90–300 K) for Cu_2I_2 ; the green line represents the fit according to the TADF equation. (g) RL spectra of the Cu_2X_2 crystals compared with reference scintillators, namely, BGO, LYSO, and the Cu_4I_4 nanocluster (50 kV and 50 μA). (h) Calculated X-ray attenuation efficiencies of Cu_2X_2 , Cu_4I_4 , DMAC-TRZ, BGO, and LYSO versus their thicknesses across the entire range of X-ray photon energies (0–50 keV) (Figure S9).

In this study, we report a series of hybrid organic–inorganic non-mechanochromic TADF Cu_2X_2 nanoclusters for X-ray scintillation and imaging with the highest light yield in their class. The TADF Cu(I) halide nanoclusters [$\text{Cu}_2\text{I}_2(\text{Dppy})(\text{F-PPh}_3)_2$],³⁵ [$\text{Cu}_2\text{Br}_2(\text{Dppy})(\text{F-PPh}_3)_2$], and [$\text{Cu}_2\text{Cl}_2(\text{Dppy})(\text{F-PPh}_3)_2$], hereafter abbreviated as Cu_2X_2 , consisting of diphenyl-2-pyridylphosphine (Dppy), tri-(4-fluorophenyl)-phosphine (F-PPh₃), and halogen atoms (I, Br, or Cl), exhibit strong photoluminescence, X-ray stopping power, and [Cu_2X_2]-to-ligand charge transfer characteristics. Consequently, these TADF nanoclusters display a record scintillation performance, having a relative light yield of up to 175 000 photons MeV^{-1} . We undertook a series of characterizations, including theoretical modeling, to elucidate the underlying mechanisms for the pronounced radioluminescence properties observed in the Cu_2X_2 nanoclusters, especially in Cu_2I_2 . Furthermore, we demonstrate that flexible Cu_2X_2 scintillation films can be fabricated by using a soft-pressing strategy. This method not only improves the quality of the scintillator but

also augments its imaging capabilities by enhancing the crystallographic alignment of the Cu_2X_2 microcrystals and reducing the optical scattering in the scintillator screens. We showcase the potential applications of these flexible scintillator films in X-ray imaging, achieving a high spatial resolution (~ 30 lp (line pairs) mm^{-1}) with a high resistance to humidity and radiation.

Single crystals of $\text{Cu}_2\text{X}_2(\text{Dppy})(\text{F-PPh}_3)_2$ ($X = \text{Cl}, \text{Br}, \text{or I}$) were obtained by the layering method with dichloromethane (DCM) and *n*-pentane.^{35,36} Notably, neither the structures of the bromide nor the chloride of this class of Cu_2X_2 nanoclusters have been previously characterized structurally. As determined by single-crystal X-ray diffraction, the Cu_2X_2 nanoclusters are isostructural and crystallize into the space group of $P2_1/c$ (Table S7).³⁵ They have similar cell parameters and butterfly shaped [Cu_2X_2] cores, which are indicative of a tetrahedral coordination geometry (Figure 1a,b and Table S7). The distance between the two copper atoms is approximately 2.70 Å in all of the obtained structures (Figure S1). The main

bond lengths and angles (X–Cu–X, P–Cu–P, and N–Cu–P) of the Cu_2X_2 nanoclusters are listed in Table S1.

The steady-state photoluminescence (PL) spectra of these Cu_2X_2 nanoclusters show broad emission bands in the visible range, centered approximately at 533 (I), 553 (Br), and 580 (Cl) nm, respectively (Figure 1d). Interestingly, a clear trend of the emission maxima is observed that is dependent on the halides, as a blue-shift occurs from Cu_2Cl_2 to Cu_2I_2 .³⁶ As shown in Figure 1c, the PL decay times at room temperature (rt) also demonstrate the same trend, gradually decreasing from 24.95 μs for Cu_2Cl_2 to 11.67 μs for Cu_2I_2 . We further investigated the PL characteristics of Cu_2X_2 by examining their temperature dependence in the range of 90–300 K. The emission energy increased as the temperature increased from 90 to 300 K, resulting in a 23, 17, and 7 nm blue-shift of the maximum emission under an excitation at 365 nm for Cu_2Cl_2 , Cu_2Br_2 , and Cu_2I_2 , respectively (Figures S2, S3, and S5). This blue-shift is a result of the thermal activation of the energetically higher-lying S_1 state above $T = 90$ K. Furthermore, this temperature-dependent enhancement in luminescence demonstrates an excellent characteristic of reverse intersystem crossing (RISC) during triplet-to-singlet up-conversion.^{37,38}

To investigate the potential TADF properties of these nanoclusters, we also measured the temperature-dependent emission decay times, which varied gradually from 90 to 300 K (Figures 1e, S2, S4, and S6). As the temperature increased, the fast RISC process to the S_1 state occurred, which initiated an additional radiative TADF process via the decay path from the S_1 state. This resulted in a noticeable blue-shift in the emission spectrum and a significant reduction in the emission decay time, as the emitting S_1 state lies higher in energy than the T_1 state. The modified Boltzmann equation was used to fit the emission decay lifetimes τ (for $T = 90$ to 300 K), which yielded an activation energy of $\Delta E(S_1-T_1) = 117$ (Cu_2Cl_2), 100 (Cu_2Br_2), and 46.5 meV (Cu_2I_2), respectively (Figures 1f, S4a, and S6a). Additionally, the radiative rates $k(S_1-S_0)$ and $k(T_1-S_0)$ that were observed for Cu_2I_2 were 1.29×10^6 and 2.89×10^4 s^{-1} , respectively, corresponding to a prompt fluorescence decay time of $\tau(S_1-S_0) = 775$ ns and a phosphorescence decay time of $\tau(T_1-S_0) = 34.6$ μs . In comparison, Cu_2Br_2 and Cu_2Cl_2 show slightly higher $k(S_1-S_0)$ and lower $k(T_1-S_0)$ values than Cu_2I_2 (Table S2). The radiative decay rate (k_r) can be determined using emission QYs and decay times, and it was found to be 2.18×10^4 (Cu_2Cl_2), 5.96×10^4 (Cu_2Br_2), and 7.78×10^4 s^{-1} (Cu_2I_2) at 300 K, respectively. At rt, the fitted curve displays a downward trend rather than a plateau, indicating that rt emission includes contributions from delayed fluorescence and phosphorescence, corresponding to triplet-state emission at low temperatures and singlet-excited-state emission at high temperatures.^{15,16,39} Additionally, the ratio $I(S_1)/I(T_1)$ at ambient temperature is approximately 2.20 (Cu_2Cl_2), 3.30 (Cu_2Br_2) and 2.50 (Cu_2I_2), respectively, indicating that TADF is the predominant factor in the emission spectra (Table S3).¹⁸ The alignment between the activation energy and the spectral shift with an increase in temperature supports the assignment of ambient temperature emission as TADF. By combining the TADF and phosphorescence pathways, the decay times of the Cu_2X_2 nanoclusters are significantly reduced compared to phosphorescent Cu(I) complexes, which typically exhibit decay times that are of the order of several hundred microseconds, even up to milliseconds.⁴⁰ In particular, Cu_2I_2 shows a smaller $\Delta E(S_1-T_1)$ and

a shorter τ_{TADF} compared to the published TADF mononuclear CuI complexes and Cu_4I_4 nanoclusters,^{9,11,25} implying stronger TADF characteristics.

Given that these inorganic–organic hybrid TADF nanoclusters have high photoluminescence QYs, low self-absorption, and a good ability to harvest both singlet and triplet excitons, we were motivated to study their potential as scintillation materials. To assess their X-ray absorption capabilities, X-ray absorptions were calculated on the Cu_2X_2 nanoclusters, and it was observed that the resonance absorption edges increased gradually from Cu_2Cl_2 to Cu_2I_2 because of the heavy atom effect (Figure S7). As shown in Figure S8, the RL spectrum showed negligible shifting compared to the PL spectrum at 300 K, indicating that the emission excited by ultraviolet (UV) and X-ray radiation likely originated from the same excited states. Respectively, the TADF Cu_2Cl_2 , Cu_2Br_2 , and Cu_2I_2 nanoclusters exhibit RL responses that are 0.86-, 0.98-, and 2.01-fold higher than the scintillation standard (LYSO ~ 33 200 photons MeV^{-1})⁴¹ with similar sample and reference thicknesses (Figure 1g and Table S4). It is important to note that the surfaces of the single-crystal reference samples were optically polished and that the Cu_2X_2 cluster samples contained microscale crystals with extremely smooth surfaces. Figure 1h shows a comparison of the X-ray attenuation efficiencies (XAEs) of the common inorganic X-ray scintillators BGO and LYSO, organic scintillator (DMAC-TRZ),¹³ Cu_4I_4 , and the Cu_2X_2 nanoclusters. The TADF Cu_2X_2 nanoclusters have relatively high XAEs, being 1 order of magnitude higher than those of organic TADF scintillating materials (Figure 1h). The Cu_2I_2 nanocluster film with a thickness of 0.5 mm absorbs 38% X-ray radiation; in comparison, one of the best organic scintillators DMAC-TRZ can only achieve such a high XAE with a thickness of 6 mm.¹³ Consequently, the light yields of Cu_2Cl_2 , Cu_2Br_2 , and Cu_2I_2 were estimated to be $98\,000 \pm 4100$, $102\,400 \pm 16\,100$, and $175\,300 \pm 4400$ photons MeV^{-1} , respectively (Figures 1g and S10 and Table S4). The method for determining the light yield is given in the Supporting Information. To further ascertain the accuracy of the light yield measurements, we compared the nanoclusters to another common scintillator, BGO, and the published nanocluster scintillator Cu_4I_4 ²⁵ under identical experimental conditions (Figures 1g and S10 and Table S4). Both standards yielded light yield values for the Cu_2X_2 nanoclusters that were in good agreement (Figure S10). Moreover, the relative light yields of Cu_2I_2 and Cu_2Br_2 were much higher than those of the reported phosphorescent tetranuclear Cu(I) iodide nanoclusters ($\text{Cu}_4\text{I}_4\text{py}_4$ and $\text{Cu}_4\text{I}_4(\text{DBA})_4$)^{11,25} and mononuclear Cu(I) halide complexes [$\text{CuI}(\text{PPh}_3)_2\text{L}$]⁹ (Table S5). These results indicate that the synergistic effect of high X-ray absorption by heavy halogen atoms (Br or I) and the harvesting of both singlet and triplet excitons through TADF behaviors leads to enhanced light yields.

To gain insight into the electronic properties and the underlying luminescence mechanism of the TADF Cu_2X_2 nanoclusters, we performed absorption spectral analysis and density functional theory (DFT) calculations. As shown in Figure S11, the absorption bands of all of the Cu_2X_2 nanoclusters are presented with extinction coefficient values and suggestions for the corresponding transitions. The electronic absorption spectra of the Cu_2X_2 nanoclusters display multiple absorption peaks in the range of 250–320 nm, which are assigned to ligand-centered (LC) $\pi-\pi^*$ transitions.^{9,42,43} In

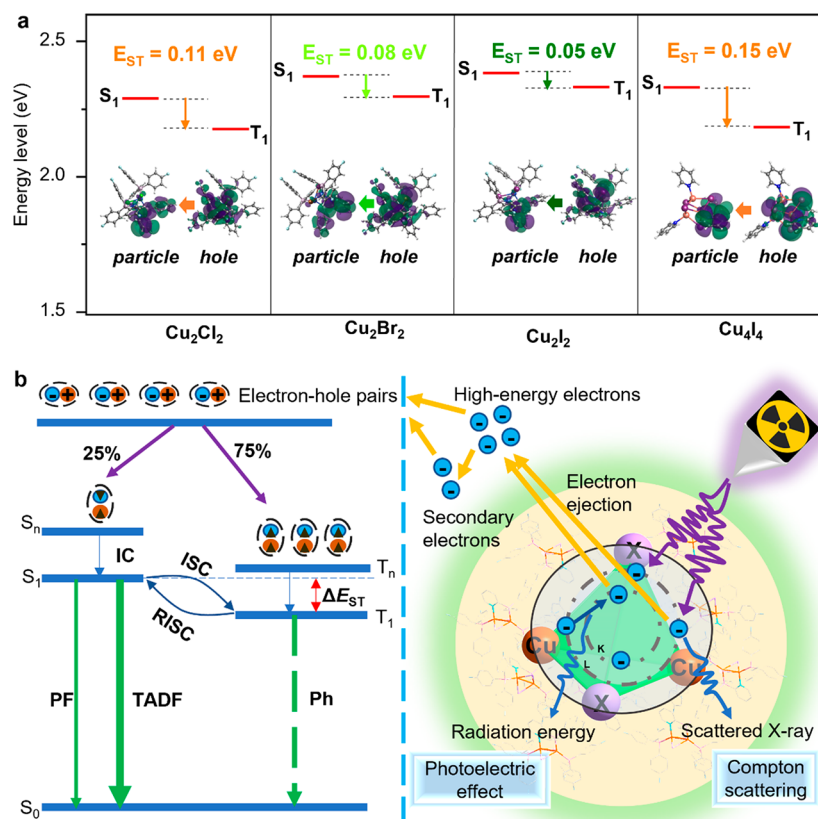


Figure 2. DFT calculations and proposed mechanism of luminescence. (a) Calculated energy levels of the lowest singlet (S_1) and triplet (T_1) states for the Cu_2X_2 and Cu_4I_4 nanoclusters. Insets show the charge densities for the holes and particles from neutral orbital transitions from the ground state (S_0) to S_1 . The DFT calculations were performed at the B3LYP/6-31G(d,p) + LANL2DZ level. (b) Proposed RL mechanism for the TADF Cu_2X_2 nanoclusters.

addition, the low-energy absorption tails in the range of 330–410 nm are assigned to (metal and halogen)-to-ligand charge transfer (M+X)LCT transitions.^{9,42,43} According to the theoretical absorption spectra of the Cu_2X_2 nanoclusters, the excited states with relatively high oscillator strength originate mainly from transitions from degenerate occupied orbitals (HOMO – 4 to HOMO (highest occupied molecular orbital)) to unoccupied orbitals ((lowest unoccupied molecular orbital) LUMO to LUMO + 1) (Figure S12 and Table S6). On the basis of the optimized geometries of S_1 and T_1 for the Cu_2X_2 and Cu_4I_4 nanoclusters, we conducted natural transition orbitals (NTOs) analysis, as shown in Figure 2a. For all three of the Cu_2X_2 nanoclusters, the optical absorption from the ground state (S_0) to S_1 can be attributed to the electronic transitions from the “hole” state delocalized in the $[\text{Cu}_2\text{X}_2]$ core to the “particle” state localized mainly within the bridging organic ligand Dppy. This suggests that the optical transition is characterized by a (M+X)LCT character between the $[\text{Cu}_2\text{X}_2]$ core and the bridging Dppy ligands, with negligible contribution from the monodentate P-donors. Structurally, the Cu–Cu interaction in Cu_2I_2 is larger than the intermolecular interactions of Cu_2Cl_2 and Cu_2Br_2 (Figure S1), which have been suggested to induce important or even unique photophysical/photochemical properties.⁴⁴ A stronger Cu–Cu interaction increases the possibility of orbital overlap, which affects luminescence performance.⁴⁴ This implies that Cu_2I_2 might emerge as one of the best TADF materials due to its potential balance between luminescence and orbital overlap. The DFT findings also point to a notable shift in the Cu_2X_2

series (Figure 2a and Table S2), with shifts approximating the $\Delta E(S_1-T_1)$ that is responsible for thermally activating TADF emission. The value of the RISC rate (k_{RISC}) is inversely correlated with ΔE_{ST} and positively correlated with the intensity of the SOC between singlet- and triplet-state orbitals.⁴⁵ Therefore, decreasing ΔE_{ST} and increasing SOC are both beneficial for increasing k_{RISC} . In addition, the heavy atom effect mainly contributes to an increase in the SOC, resulting in a decrease in ΔE_{ST} . Therefore, this effect results in a higher radiative TADF decay rate, allowing the emission QY to reach 86% for Cu_2I_2 at ambient temperature. Also, the emission maximum sequence (580 (Cu_2Cl_2), 550 (Cu_2Br_2), and 533 nm (Cu_2I_2)) corresponds to the order of the halide ligand field strengths ($\text{I}^- < \text{Br}^- < \text{Cl}^-$), leading to stronger d-orbital splitting and HOMO destabilization, thus ultimately reducing the energy gap. The results reveal a gradual decrease in the potential energy level difference due to the introduction of Br and I, significantly enhancing the orbital coupling between the singlet and triplet states. According to the above results, Cu_2I_2 has the strongest TADF character among the Cu_2X_2 complexes (Table S2). Figures S13 and S14 shows the calculated spin-orbit coupling matrix element (SOCME) values and oscillator strengths in the optimized S_1 of the Cu_2I_2 and Cu_4I_4 nanoclusters. Compared to the well-known cubane-like Cu_4I_4 nanoclusters, Cu_2I_2 shows a smaller ΔE_{ST} and a higher SOC between the singlet- and triplet-state orbitals, which are both beneficial for increasing k_{RISC} and greatly improving the ability to convert ionizing radiation into low-energy photons.^{45,46}

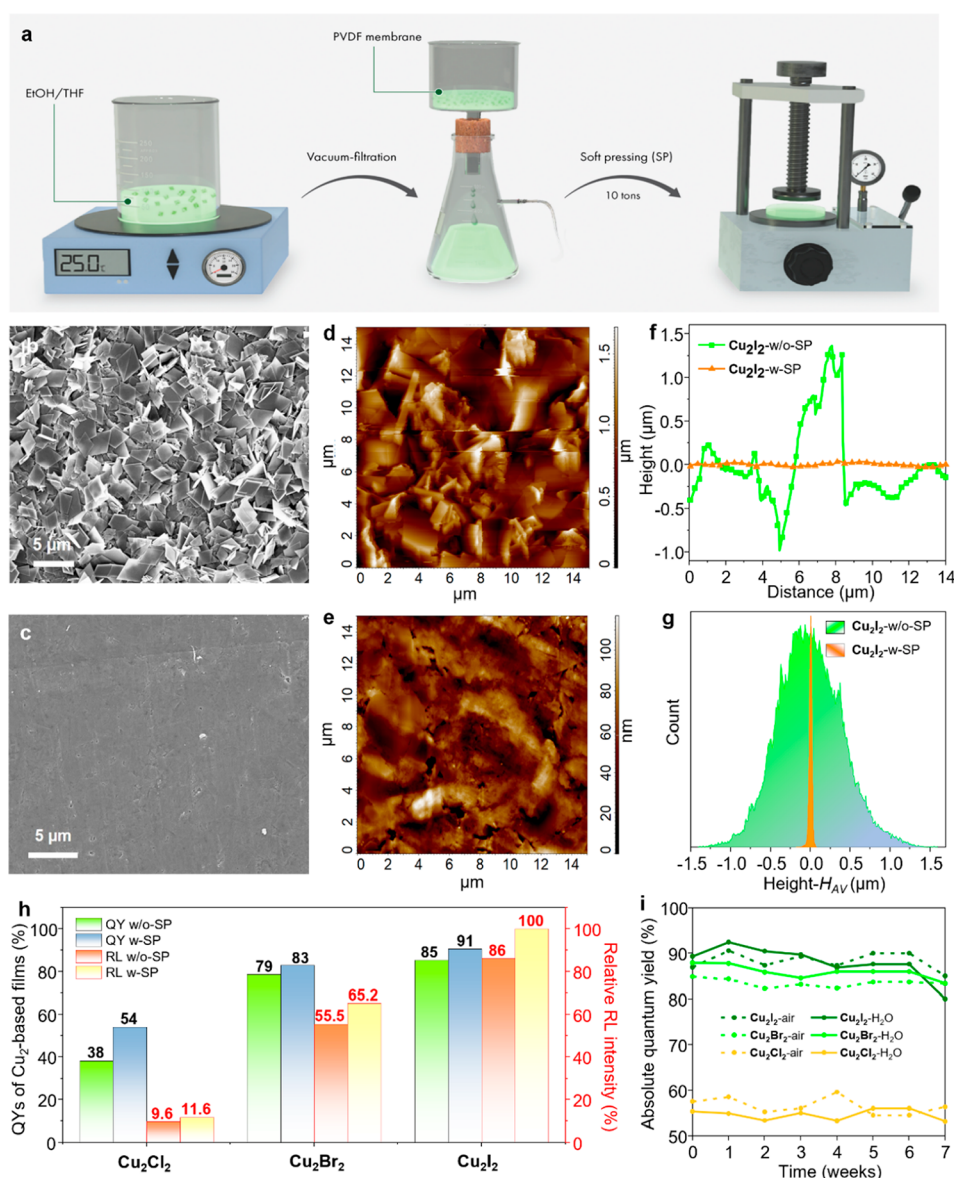


Figure 3. Characterization of the Cu₂X₂ films. (a) Scheme of the Cu₂X₂ film synthesis procedure by suction filtration (w/o-SP) and soft pressing (w-SP). SEM images of the front view of (b) Cu₂I₂-w/o-SP and (c) Cu₂I₂-w-SP. AFM images of (d) Cu₂I₂-w/o-SP and (e) Cu₂I₂-w-SP. (f) AFM line segment and (g) AFM height distribution analyses. (h) Comparison of QY and RL intensities of the Cu₂X₂ films (the integrated RL intensity of Cu₂I₂-w-SP was normalized to 100%). (i) The stability of these Cu₂X₂-w-SP films in an ambient atmosphere and in water.

To understand the nature of the excited state of luminescent materials under high-energy X-ray radiation and low-energy photoexcitation, the PL and RL spectra of Cu₂X₂ were compared at $T = 300$ K (Figure S8). The agreement of both spectra indicates that, whether excited by UV or X-ray photons, the emission likely originates from the same excited states. Figure 2b schematically illustrates the proposed luminescence mechanism under X-ray irradiation for the Cu₂X₂ nanoclusters. During the process of X-ray absorption and conversion, X-rays primarily interact with the heavy [Cu₂X₂] cores (Cu and X, especially Br and I) through the photoelectric effect and Compton scattering mechanism, resulting in the release of energetic electrons.^{13,47} After being generated by electron–electron scattering and Auger processes, these hot electrons undergo rapid thermalization and are subsequently captured by the organic ligands within the

Cu₂X₂ nanocluster scintillators. Such interactions result in the creation of hole–electron pairs and a subsequent cascade of secondary electrons that exhibit characteristic (M+X)LCT and LCT states.^{9,42} Thereafter, ~25% of singlet excitons and ~75% of triplet excitons are produced by charge recombination between the electrons and holes.^{15,21} Finally, given the minimal ΔE_{ST} between S₁ and T₁ of the Cu₂X₂ nanoclusters, the triplet excitons can efficiently transfer to singlet excitons via a RISC process. This results in efficient (M+X)LCT and LCT states for radiative processes, boasting higher internal quantum efficiencies and light yields. The highly radiative (M+X)LCT and LCT states, combined with the excellent RISC process, play a crucial role in enhancing the RL in the Cu₂X₂ nanoclusters, thereby making them promising candidates for X-ray imaging.⁴²

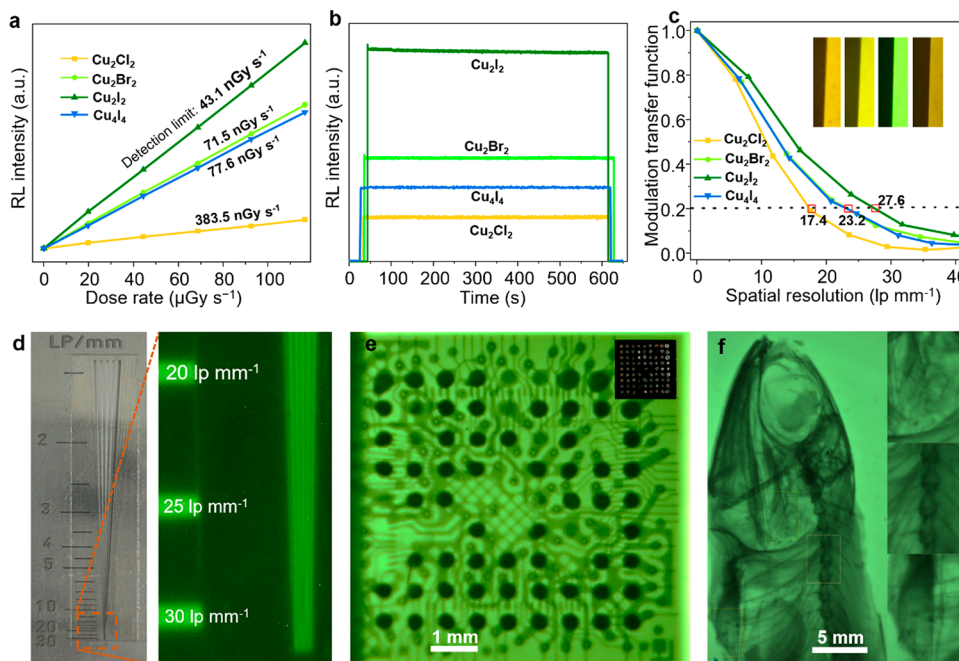


Figure 4. X-ray imaging with the copper halide nanoclusters. (a) Detection limits of the Cu_2 -based TADF films and the Cu_4I_4 film. (b) Normalized RL intensity at the corresponding emission maxima of the Cu_2X_2 -w-SP TADF films and Cu_4I_4 film under continuous X-ray irradiation (dose rate of $146.5 \mu\text{Gy s}^{-1}$). (c) Modulation transfer function (MTF) curves of the Cu_2X_2 -w-SP TADF films and Cu_4I_4 film measured by the slanted-edge method (inset: corresponding X-ray edge images of the Cu_2X_2 -w-SP TADF films). Bright-field and X-ray images using the Cu_2I_2 -w-SP TADF film of (d) a standard line-pair card, (e) an electronic chip, and (f) a small fish (dose rate of $130 \mu\text{Gy/s}$; exposure of 10 s).

In recent years, there has been a growing interest in using environmentally friendly Cu-based materials with high luminescence efficiencies as scintillators.^{7,9,11,25,41,48} The properties of these materials, however, are highly dependent on their manufacturing strategy. The most widely used strategies include polymer-based flexible films,^{9,25,26,42} high-temperature transparent ceramics,^{32,49} and high-pressure wafers.^{11,33} Despite their promising properties, most Cu-based compounds exhibit strong mechanochromic luminescence^{30,31} and aggregation-induced luminescence,^{50,51} making it challenging to optimize the luminescence efficiency with respect to thickness and size using current preparation strategies. Fortunately, the recent development of a vacuum filtration strategy has provided a pathway to address these problems and has yielded excellent results for the application of nanoscale materials.^{6,52}

A vacuum filtration strategy was successfully applied to the TADF Cu_2X_2 nanoclusters for the preparation of high-quality scintillator screens (Figure 3a). Because of their excellent solubility in tetrahydrofuran (THF), uniform rhomboidal microcrystals were first achieved by using ethanol as an antisolvent (the reaction is shown in the Supporting Information). The microcrystals, which had a relatively uniform size of about $2 \mu\text{m}$, were analyzed by transmission electron microscopy (TEM) (Figures S15–S17) to confirm the elemental distributions. The obtained Cu_2X_2 microcrystals were then uniformly dispersed in a mixed solvent of THF and ethanol for the fabrication of stable and flexible screens. For vacuum filtration, a hydrophilic poly(vinylidene fluoride) (PVDF) filter membrane (Figure S18) was used to trap and support the hydrophobic Cu_2X_2 microcrystals. Scanning electron microscopy (SEM) was used to characterize the surface of the Cu_2X_2 films fabricated by vacuum filtration

(abbreviated as Cu_2X_2 -w/o-SP). The thicknesses of the Cu_2X_2 microcrystal layer and the PVDF membrane layer after vacuum filtration were 50 and $80 \mu\text{m}$, respectively (Figure S19a). As shown in Figure 3b, the disordered surface arrangement of the PVDF membrane resulted in an uneven surface with numerous defects between adjacent Cu_2X_2 microcrystals, which could negatively impact future applications. To address this issue, soft-pressing was subsequently employed to obtain uniform and smooth films of the Cu_2X_2 microcrystals at a positive pressure of 10 tons (Figure 3a). The Cu_2X_2 -based films after being soft-pressed (abbreviated as Cu_2X_2 -w-SP) had a compact structure and a relatively flat surface (Figures 3c, S20, and S21), and the thicknesses of the Cu_2X_2 microcrystal layer and the PVDF membrane layer were reduced to 20 and $60 \mu\text{m}$, respectively (Figure S19b). The surface quality of the achieved films was further evaluated by using atomic force microscopy (AFM) (Figure 3d,e). The root-mean-square roughnesses of the Cu_2X_2 -w/o-SP and Cu_2X_2 -w-SP films were calculated to be 400 and 19 nm, respectively, for areas of $15 \times 15 \mu\text{m}^2$. The roughness of the Cu_2X_2 -w-SP film was greatly reduced compared to that of the Cu_2X_2 -w/o-SP film, as evidenced by both the AFM line segment and height distribution analyses (panels f and g of Figure 3, respectively). Clearly, soft-pressing using 10 tons of pressure was sufficient in eliminating pores in the thick film and rearranging the Cu_2X_2 microcrystals. The obtained Cu_2X_2 -w-SP films exhibited excellent uniformity under both UV and X-ray radiation, and the thickness and flexibility could be controlled by the amount of Cu_2X_2 microcrystals in the films (Figure S22).

Furthermore, the pressure stability greatly affects the subsequent scintillator properties of the Cu_2X_2 -based films. Unlike most Cu–I nanoclusters (especially Cu_4I_4), which have excellent mechanochromic properties due to their structural

variation (Figure S23),^{38,53} the good pressure stability of Cu_2X_2 was confirmed by powder X-ray diffraction (PXRD) diagrams and emission spectra before and after soft-pressing (Figure S24). In addition, the intensity of the (002) peak was found to be significantly enhanced, especially for Cu_2I_2 (Figure S24c), indicating a preferred orientation after laminating. Moreover, all of the Cu_2X_2 complexes exhibited large Stokes shifts, almost negligible self-adsorption, and high photoluminescence quantum yields (PLQYs) at room temperature (54% for Cu_2Cl_2 -w-SP, 83% for Cu_2Br_2 -w-SP, and 91% for Cu_2I_2 -w-SP). The decrease in the amount of voids and pinholes in the soft-pressed Cu_2X_2 -based films resulted in a remarkable enhancement in film quality, as evidenced by the increasing intensity of PL and RL, as shown in Figures 3h and S25. The substantial reduction in the porosity of the films effectively hinders ion migration, minimizes light-induced carrier trapping and recombination, and consequently leads to a reduction in X-ray scattering. Additionally, the Cu_2X_2 -w-SP films showed excellent stability in both ambient atmosphere and water, surpassing traditional bulk X-ray scintillating and metal halide perovskites^{47,53} (Figures 3i, S26, and S27).

To verify the scintillator properties of the Cu_2X_2 -w-SP screens, linear responses were first obtained at higher dose rates ranging from 19.8 to 116.4 $\mu\text{Gy s}^{-1}$ (Figure S28), making them an excellent choice for achieving a good X-ray imaging contrast. The limit of detection (LOD) was determined based on a signal-to-noise ratio (SNR) of 3. In terms of response sensitivity, the inclusion of heavy elements (I and Br) significantly enhanced the response sensitivity of the scintillators to X-ray dose rates, with Cu_2I_2 -w-SP exhibiting a response sensitivity of 43.1 nGy s^{-1} , as opposed to those of 383.5 and 71.5 nGy s^{-1} for Cu_2Cl_2 -w-SP and Cu_2Br_2 -w-SP (almost the same as those of Cu_4I_4), respectively (Figure 4a). Remarkably, the X-ray dose rate of Cu_2I_2 -w-SP was approximately 100 times lower than the requirement for standard diagnostic imaging (5.5 $\mu\text{Gy s}^{-1}$). Additionally, all of the Cu_2X_2 -w-SP screens exhibited excellent RL stability, with the RL intensity remaining above 90% even when exposed to ionizing radiation at a dose rate of 146.5 $\mu\text{Gy s}^{-1}$ for a continuous duration of 600 s (Figure 4b).

Motivated by the promising luminescence properties of the TADF Cu_2X_2 -w-SP screens, we investigated their imaging resolutions. The imaging resolutions followed the trend of RL intensity, with high resolutions being achieved for the Cu_2X_2 -w-SP screen (Figures 4d and S29), according to a standard line-pair card. By calculating the modulation transfer function (MTF) of the standard X-ray slant image, we determined the spatial resolutions (at an MTF of 0.2; see Figures 4c and S30–S32) to be 17.4, 23.2, and 27.6 lp mm^{-1} for Cu_2Cl_2 , Cu_2Br_2 , and Cu_2I_2 , respectively. The high spatial imaging resolution of the TADF Cu_2I_2 -w-SP screen surpasses that of most reported organic and inorganic materials, e.g., the Cu_4I_4 nanocluster ($\sim 23 \text{ lp mm}^{-1}$), further underscoring the promising practical potential of incorporating TADF chromophores into high-performance X-ray imaging scintillators using pressing engineering strategies. To demonstrate the practical value of the TADF scintillators, we conducted a series of imaging experiments. In particular, we employed the principle of X-ray contrast imaging to inspect the complex internal structures of two electronic chips of different sizes (Figures 4e, S33, and S34), which were opaque to visible light. The Cu_2X_2 -w-SP scintillator screens enabled a clear visualization of the complex structure of the electronic chip. Figures 4e and S33 display the

X-ray imaging of a targeted chip positioned on the complex Cu_2I_2 -w-SP scintillator screen. Owing to the differential absorption of X-rays by different materials in the chip, the internal structure and circuitry beneath the black plate are clearly visible. A small sample of fish was also used to evaluate the potential of the Cu_2X_2 -w-SP scintillator screens for biological applications (Figure 4f). As shown in Figures 4f and S34, the skeleton, head, and fin of the fish are clearly observable under X-ray exposure. We believe that such high resolution and superior imaging capability of the Cu_2I_2 -w-SP TADF scintillator screen exceed those reported for most scintillators.^{8–10,12,13,25,49,54,55} These results further demonstrate the promising potential of organic–inorganic hybrid TADF scintillators in medical diagnosis and industrial inspection.

In this study, we demonstrated a series of TADF Cu_2X_2 nanoclusters, which exhibit intense radioluminescence, low self-absorption, and high radiation stability. The efficient [Cu_2X_2]-to-ligand charge transfer characteristics of the nanoclusters contribute to their scintillation performance, which is superior to that of other copper halide nanoclusters, resulting in record-high relative light yields (as high as 175 000 photons MeV^{-1} for Cu_2I_2). Because of their non-mechanochromic properties, flexible Cu_2X_2 TADF scintillation screens were successfully created through a soft-pressing strategy. These fabricated screens showcased an outstanding performance, having a detection limit as low as 43.1 nGy s^{-1} , an X-ray imaging resolution as high as 27.6 lp mm^{-1} , and a highly linear response to the X-ray dose rate. These findings highlight the potential of inorganic–organic hybrid metal nanoclusters in advancing the development of high-energy radiation scintillation materials.

■ ASSOCIATED CONTENT

Supporting Information

The Supporting Information is available free of charge at <https://pubs.acs.org/doi/10.1021/acsenerylett.3c02050>.

Experimental section (material synthesis, material characterization descriptions, time-dependent DFT calculations, and light yield, XAE, and X-ray imaging spatial resolution calculations), temperature-dependent emission intensities, temperature-dependent decay time measurements, absorption spectra calculations, TEM and SEM images of crystals and scintillating films, EDS mapping of crystals, XRD analysis, quantum yield determination, X-ray images of objects, comparisons of light yields, comparison of X-ray detection and imaging with TADF scintillators, crystal data for the Cu_2X_2 nanoclusters, and details of literature comparisons (PDF)

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Notes

The authors declare the following competing financial interest(s): O.M.B. is a founder of Quantum Solutions, a company that develops optoelectronic devices.

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